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## LETTER TO THE EDITOR

## The Lie derivative of the vorticity vector in an isometric flow

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#### Abstract

It is shown that in an isometric flow the vorticity vector is invariant under the group of motions.


A fluid flow in general relativity is said to be isometric if there exists a time-like Killing vector $\xi_{\alpha}$ colinear with the fluid velocity vector $u_{\alpha}$, ie if

$$
\begin{equation*}
u_{\alpha}=\xi_{\alpha} / \xi, \quad \xi^{2}=\xi_{\alpha} \xi^{\alpha}>0, \quad \xi_{\alpha ; \beta}+\xi_{\beta ; \alpha}=0 \tag{1}
\end{equation*}
$$

Ciubotariu (1972a, b) has considered such flows in detail and has shown that

$$
\begin{align*}
& \mathscr{L}(\xi) \omega_{\alpha \beta}=0  \tag{2}\\
& h_{\alpha} \mathscr{L}(u) \omega^{\alpha}=0 \tag{3}
\end{align*}
$$

where $\mathscr{L}(\xi)$ and $\mathscr{L}(u)$ denote the Lie derivatives with respect to $\xi_{\alpha}$ and $u_{\alpha}$ respectively, $\omega_{\alpha \beta}$ and $\omega^{\alpha}$ are the vorticity tensor and the vorticity vector and $h_{\alpha}$ is the magnetic field four-vector. Result (3) is important as it can be regarded as a generalization of Ferraro's (1937) theorem of isorotation. It is the purpose of this note to show that besides (2) and (3) we also have in an isometric flow

$$
\begin{equation*}
\mathscr{L}(\xi) \omega^{\alpha}=0 \tag{4}
\end{equation*}
$$

Instead of starting from first principles, we make use of (2). The vorticity vector $\omega^{x}$ can be expressed in terms of $\omega_{\alpha \beta}$ as

$$
\begin{equation*}
\omega^{\alpha}=\frac{1}{2} \eta^{\alpha \beta \lambda \mu} \omega_{\beta \lambda} u_{\mu} \tag{5}
\end{equation*}
$$

where $\eta^{\alpha \beta \lambda \mu}$ is the permutation tensor. As

$$
\begin{equation*}
\eta_{\alpha \sigma \rho v} \eta^{\alpha \beta \lambda \mu}=-\delta_{\sigma}^{\beta}\left(\delta_{\rho}^{\lambda} \delta_{v}^{\mu}-\delta_{\rho}^{\mu} \delta_{v}^{\lambda}\right)-\delta_{\sigma}^{\lambda}\left(\delta_{\rho}^{\mu} \delta_{v}^{\beta}-\delta_{\rho}^{\beta} \delta_{v}^{\mu}\right)-\delta_{\sigma}^{\mu}\left(\delta_{\rho}^{\beta} \delta_{v}^{\lambda}-\delta_{\rho}^{\lambda} \delta_{v}^{\beta}\right) \tag{6}
\end{equation*}
$$

we have

$$
\begin{equation*}
\eta_{\sigma \rho v x} \omega^{\alpha}=\omega_{\sigma \rho} u_{v}+\omega_{\rho v} u_{\sigma}+\omega_{v \sigma} u_{\rho} \tag{7}
\end{equation*}
$$

Now as $\xi_{\alpha}$ is a Killing vector,

$$
\begin{equation*}
\mathscr{L}(\xi) u_{\alpha}=0 \tag{8}
\end{equation*}
$$

Thus on taking the Lie derivative with respect to $\xi_{\alpha}$ of both sides of (7), we have

$$
\begin{equation*}
\eta_{\sigma \rho v \alpha} \mathscr{L}(\xi) \omega^{\alpha}=-\omega^{\alpha} \mathscr{L}(\xi) \eta_{\sigma \rho v \alpha} \tag{9}
\end{equation*}
$$

If we write $\mathscr{L}(\xi) \eta_{\text {ovvz }}$ out explicitly and operate on both sides of (9) by $\eta^{\sigma \nu \vee \beta}$ and note that

$$
\begin{equation*}
\eta^{\sigma \nu \vee \beta} \eta_{\sigma \rho v \alpha}=-6 \delta_{\alpha}^{\beta} \tag{10}
\end{equation*}
$$

and, for example, that

$$
\begin{equation*}
\eta^{\sigma \nu \nu \beta} \eta_{\sigma \rho \tau z}=-2\left(\delta_{\eta}^{\nu} \delta_{\alpha}^{\beta}-\delta_{\tau}^{\beta} \delta_{a}^{\nu}\right) \tag{11}
\end{equation*}
$$

then we obtain

$$
\begin{equation*}
\mathscr{L}(\xi) \omega^{\beta}=-\omega^{\beta} \xi_{; \tau} . \tag{12}
\end{equation*}
$$

But as $\xi_{\alpha}$ is a Killing vector, $\xi_{; ~}^{\tau}=0$, so that

$$
\begin{equation*}
\mathscr{L}(\xi) \omega^{\beta}=0 . \tag{13}
\end{equation*}
$$

This is the same as saying that in an isometric flow, $\omega^{\beta}$ is invariant under the group of motions.

## References

